

TRANSP physics and application to JET plasmas

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TRANSP is a time-dependent, 1 ½ D tool for interpretive and predictive analysis of tokamak, ST and RFP plasmas



Output of TRANSP (Plasma State File) is standardized for simplifying input to other computationally intensive codes

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TRANSP documentation



- User guide to select the physics modules and options is available, but no manual describing physics
- Information on TRANSP physics presented here comes from publications and discussions with TRANSP team during last 10 years
- This talk includes physics description when available and the options for physics modules in TRANSP
- Ref. to TRANSP: R. J. Goldston et al., J. Comput. Phys. 43, 61 (1981). Refs. to TRANSP modules on JET/TRANSP and NTCC pages

Outline



- Equilibrium (see also talk by Hyun-Tae Kim)
- Diagnostics simulations and data consistency: will be addressed later on
- Poloidal field diffusion
- Auxiliary heating
- Edge particle source
- MHD: sawtooth model
- Predictive TRANSP

Grad-Shafranov equation



$$\Delta^* \psi = r^2 \mathrm{div} \frac{\nabla \psi}{r^2} = -4\pi^2 \left(\mu_0 r^2 \frac{\partial p}{\partial \psi} + I \frac{\partial I}{\partial \psi} \right).$$

$$\Delta^* \psi = \frac{2\pi\mu_0 R_0 J}{\langle B^2/B_0^2 \rangle} j_{\parallel} + \frac{2\pi\mu_0 R_0^2}{B_0 \rho \mu} \left(\frac{J^2}{\langle B^2/B_0^2 \rangle} - \frac{r^2}{R_0^2} \right) \frac{\partial p}{\partial \rho}$$
$$= 2\pi\mu_0 R_0 \left[\frac{J}{\langle B^2/B_0^2 \rangle} \left(j_{\parallel} + \frac{R_0 J}{B_0 \rho \mu} \frac{\partial p}{\partial \rho} \right) - \frac{r^2}{B_0 R_0 \rho \mu} \frac{\partial p}{\partial \rho} \right]$$

- No explicit time dependence: plasma is always in equilibrium, fast relaxation process with respect to transport
- Coupled to time-dependent transport equations in predictive part of code

Equilibrium modules: VMEC



VMEC = Variational Moments Equilibrium Code (S. Hirshman, ORNL)

- the full 3D MHD equilibrium equations for arbitrary geometry
- truncated to a 2D code suitable for modeling tokamak geometries of arbitrary moment and adapted it to TRANSP:

1. fixed boundary: the plasma boundary prescribed by a set of Fourier coefficients in R and Z $\,$

- 2. other input parameters:
- a) enclosed toroidal flux
- b) pressure profile
- c) profile $\mu = \partial \psi / \partial \Phi$

3. can handle pressure anisotropies IN PRINCIPLE. The version currently in TRANSP, however, is purely isotropic

4. Φ_{tot} is used as an initial guess in arriving at a solution that conserves Ip, by varying Φ_{tot}

References:

- S.P.Hirshman and J.C.Whitson, PHYS.FLUIDS 26, 3553 (1983).
- S.P.Hirshman and H.K.Meier, PHYS.FLUIDS 28, 1387 (1985).
- S.P.Hirshman and D.K.Lee, COMP.PHYS.COMM. 39, 161 (1986).

Equilibrium modules: TEQ



- MHD equilibrium code used in the Corsica transport code (LLNL)
- Fixed boundary solution using the pressure and q profiles as input
- The vacuum R*Btor is used as a boundary condition
- After the initial startup, TEQ is called in a loop which adjusts the q profile near the edge region in order to match to the total plasma current
- Radial grid: uniform or stretched near the axis or the edge

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- Prescribed q-profile, q(time, r)
- Evolve the q profile using input data: Bpol/Btor vs (R,t) or θ vs (R,t) where tan(θ)=Bpol/Btor

Poloidal field diffusion equation (PFDE) can be used to estimate the resistivity profile in these two cases, but non-physical negative values of resistivity can come out (it depends on the quality of the q profile data, dq/dt, etc)

• Solve poloidal field diffusion equation

It is possible to switch back and forth amongst these options in the course of a run.

Poloidal field diffusion equation (PFDE)



$$\begin{split} \sigma_{\parallel} \left(\frac{\partial \psi}{\partial t} - \frac{\rho \dot{B}_{0}}{2B_{0}} \frac{\partial \psi}{\partial \rho} \right) &= \frac{J^{2} R_{0}}{\mu_{0} \rho} \frac{\partial}{\partial \rho} \left(\frac{G_{2}}{J} \frac{\partial \psi}{\partial \rho} \right) - \frac{V'}{2\pi \rho} \left(j_{BS} + j_{CD} \right) \\ G_{2} \stackrel{\text{def}}{=} \frac{V'}{4\pi^{2}} \left\langle \left(\frac{\nabla \rho}{r} \right)^{2} \right\rangle \qquad J \stackrel{\text{def}}{=} \frac{I}{R_{0} B_{0}}. \qquad \rho \stackrel{\text{def}}{=} \sqrt{\frac{\Phi}{\pi B_{0}}}. \qquad V' = \frac{\partial V}{\partial \rho}, \\ I_{\text{pl}} \stackrel{\text{def}}{=} \int_{S_{\zeta}} \mathbf{j} \cdot d\mathbf{S}_{\zeta} = \frac{1}{2\pi} \int_{V} \left(\mathbf{j} \cdot \nabla \zeta \right) d^{3}x = \frac{1}{2\pi R_{0}} \int_{0}^{\rho} V' j_{\text{tor}} d\rho = \frac{J}{2\pi R_{0}} \int_{0}^{\rho} \frac{V'}{J^{2}} j_{\parallel} d\rho \\ &= \frac{G_{2}}{\mu_{0}} \frac{\partial \psi}{\partial \rho} = \frac{2\pi B_{0}}{\mu_{0}} \rho G_{2} \mu = \frac{2\pi R_{0}}{\mu_{0}} G_{2} B_{p} \\ j_{\text{tor}} = 2\pi R_{0} \frac{\partial I_{\text{pl}}}{\partial V}, \qquad j_{\parallel} = 2\pi R_{0} J^{2} \frac{\partial}{\partial V} \left(\frac{I_{\text{pl}}}{J} \right) = J^{2} \frac{\partial}{\partial V} \int_{0}^{V} \frac{j_{\text{tor}}}{J} dV. \end{split}$$

$$j_{\parallel} = \sigma_{\parallel} E_{\parallel} + j_{\rm BS} + j_{\rm CD}$$

Initial condition for PFDE



$$\left.\psi(
ho,t)
ight|_{t=0}=\psi_0(
ho)$$
 not measured

Measured /reconstructed variables are

$$q \stackrel{\text{def}}{=} \frac{1}{\mu} = \frac{\partial \Phi}{\partial \psi}.$$

$$j_{\parallel} \stackrel{\text{def}}{=} \frac{\langle \mathbf{j} \cdot \mathbf{B} \rangle}{B_0} = \frac{2\pi R_0}{\mu_0 V'} J^2 \frac{\partial}{\partial \rho} \left(G_2 J^{-1} \frac{\partial \psi}{\partial \rho} \right)$$

but they involve 1st or 2nd derivative of $\psi(\rho)$ and equilibrium (V', G₂)

Two types of initial conditions:



Boundary conditions for PFDE



$$\left. \frac{\partial \psi}{\partial \rho} \right|_{\rho=0} = 0. ??$$

a) Prescribed total plasma current:

$$\left. \frac{\partial \psi}{\partial \rho} \right|_{\rho = \rho_B} = \frac{\mu_0}{G_2(\rho_B)} I_{\rm pl}(t)$$

b) Prescribed loop voltage: may be inconsistent with total plasma current



- Spitzer resistivity
- Analytical neoclassical models (S. P. Hirshman, et al Nucl. Fusion 1977; TSC model)
- NCLASS (full multi-species representation of plasma profiles, valid for arbitrary geometry and collisionality regimes) [W. A. Houlberg et al, Phys. Plasmas 1997]
- Sauter model (analytical expressions fitting the code simulations: Fokker–Planck equation with full collision operator, arbitrary equilibrium and collisionality regimes) [O. Sauter and C. Angioni, PoP 1999]



lbs, simulated q, NCLASS	0.36 MA
lbs, simulated q, Sauter model	0.43 MA
lbs, EFIT/q, NCLASS	0.4 MA
lbs, EFIT/q, Sauter model	0.5 MA

Current diffusion studies for JET: OH current ramp up

- Simulations of current diffusion with measured T_e and line averaged Z_{eff}
- EFIT/Q is initial profile
- First sawtooth crash (i.e. q₀ > 0.8) at 6.9 s

 $\sigma \sim T_e^{3/2} / Z_{eff}$

- Sensitivity to Z_{eff} profile:
- flat Z_{eff} → slower current diffusion towards the centre, higher q₀ at the beginning of the ramp up
- peaked $Z_{eff} \rightarrow$ broader stationary j, higher q_0
- Sensitivity to T_e profile:
- peaked T_e → slower current diffusion towards the centre, higher q₀ at the beginning of the ramp up
- flat $T_e \rightarrow$ broader stationary j, higher q_0





Current diffusion studies for JET: high β_N scenario



- Current diffusion simulations with measured Te, Zeff. NCLASS is used
- EFIT/Q is initial profile taken 0.5 s before the NBI start
- Good agreement with EFIT for discharges with early and late NBI start
- Over-estimated central q in discharge with strong n=1 mode





Voitsekhovitch et al, Nucl. Fusion 2009

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Neutral beam injection: NUBEAM [A.Pankin, CPC]



- 2D code, realistic geometry
- Monte-Carlo technique for beam ions and fusion products
- FLR effects included
- Deposition, secondary CX, slowing down
- Orbit losses, collisional and anomalous diffusion
- Sawtooth mixing of fast ions
- Built into JET database
- Frequently used output:
- heat, particle and momentum source
- Fast particle distribution function $f(\rho, pol.angle, E, pitch angle)$



Examples of NUBEAM simulations for JET



NBI heating profile obtained with different beam configuration and plasma density





Voitsekhovitch et al, Nucl. Fusion 2007



Deuterium beam density profile in high density H-mode plasmas, obtained by integrating the 4D distribution function

Ion Cyclotron Heating: TORIC



- FLR full wave code
- Solves Maxwell's equations in presence of plasma and wave antenna
- Retains the 2nd harmonic wave frequency
- specify both the damping power density on minority fast ions, and the 2D wave field (E+, polarization, k_⊥, k_{||})
- $Im(E_+),$ $Re(E_+),$ V/m V/m 200 200 112.00 108 84.00 76 100 10056:00 44 28.00 12 0.14-20 -28.00-52 -56.00-100-100-84 -84.00-116-112.00-200-200-140.00
- Combined with FP module
 FPPMOD: re-normalises the original
 QL operator zone by zone while
 keeping the total power constant

Application for:

(1) minority heating (H and He3)

-150 -100 -50

50 100

- (2) fundamental D heating
- (3) mode conversion

<u>Refs</u>: M. Brambilla, PPCF 41, 1, (1999) & M. Brambilla and T. Krucken, NF 28, 1813 (1988); D. G. Swanson, Phys. Fluids 24, 2035 (1981); P. T. Colestock and R. J. Kashuba, NF 23 763 (1983); J. C. Wright et al, PoP 11, 2473 (2004)

-150 -100 -50

0 50 100

$E_{+}=\exp(i\tau)\{E_{\psi}+i(\cos\theta E_{\tau}-\sin\theta E_{\phi})\}/2^{0.5} \qquad J. \text{ Conboy}$

Lower hybrid heating and current drive



LSC (D.W.Ignat, E.J.Valeo):

- multiple ray tracing in general non-circular axisymmetric plasmas specified by a numerical equilibrium
- solve at each of several flux surfaces a simple one-dimensional (in velocity) Fokker Planck equation for the quasilinear evolution of the electron distribution function
- provide the RF power and RF-driven current electric field

Refs: D. W. Ignat, Phys. Fluids, 1981; E. J. Valeo and D. C. Eder, J. Comp. Physics, 1987. C. F. F. Karney and N. J. Fisch, Phys. Fluids 1986

GENRAY (work in progress):

- ray tracing code

- can be used for modeling electron cyclotron and lower hybrid heating and current drive in tokamaks

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- Fluid 1.5D model, no poloidal variation of the source
- Simplified geometry: nested cylinders of shifted centers given by flux surfaces
- Multiple ion and neutral species
- Charge exchange, impact ionization
- Input: neutral flux and energy of incoming neutrals (multiple energy species)
- For JET simulations: gas puff and renormalised Dα signal, Γneut = Γpuff + const*Γα
- Wall particle source Γneut is ad hoc, should not be used for estimation of global particle confinement
- TRANSP-EDGE2D estimation of wall particle source for JET plasmas started under ITM/ISM (collaboration with Paula Belo)

TRANS-EDGE2D iterations (more detailed description in Voitsekhovitch et al NF 2013)



Estimation of particle confinement time and Γα normalisation [Voitsekhovitch et al EPS 2012: 79635 (low power HS): τp=0.4 s, τE=0.16 s, const=17.5 77922 (high power HS): τp=0.54 s, τE=0.25 s, const=16.8 Effect of gas puff (CX losses) on neutron yield during current ramp up



Γpuff = C*GASM, **C=1**, 20, 50



sensitivity to CX losses is tested by varying gas puff

- good agreement with measured neutron yield with C ~ 40, but the CX losses are high

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Kadomtsev reconnection model [Sov. J. Plasma Phys., 1975]

TRANSP sawtooth models:

- crash times taken from data (ECE, soft X-ray)
- radius with q=1 is extended to provide the helical flux conservation
- q, Te, Ti, thermal and fast ion (NBI & RF) densities are mixed following the reconnection of magnetic flux surfaces:
 - → centrally peaked profiles flatten
 - \rightarrow off-axis particles go towards the center









Kadomtsev mixing yields q>=1 everywhere; Porcelli mixing generally leaves a region with q<1 near the axis

Options in Porcelli model:

- two mixing regions: "island" around q=1 and axial region inside the island annulus
- single mixing region for predicted plasma species, covering both the q=1 island and the axial region
- island width controlled by user

- fraction of mixed fast ions controlled by user

Evolution of q-profile during the first sawtooth crash in Kadomtsev (red and pink) and Porcelli (red and blue) models (same initial conditions)



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Predictive TRANSP:

- thermal ion and electron temperature equations
- transport modules for temperature and density
- momentum equation
- modelling of density of trace species



Particle:
$$\frac{\partial n}{\partial t} = -\frac{1}{r} \frac{\partial}{\partial r} (rnV) + S_{gas} + S_{beam} - L_{c-x, recomb, etc}$$

Energy (e⁻):
$$\frac{\partial nT}{\partial t} = -\frac{1}{r}\frac{\partial}{\partial r}\left(rn\chi\frac{\partial}{\partial r}T\right) - nTV_r + P_{OH} + P_{beam} - P_{ie} - P_{rad}$$

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- Whole plasma region or shifted boundary
- Source computed with NUBEAM, TORIC, LSC, FRANTIC
- Radiative power taken from measurements or computed assuming coronal equilibrium
- Transport models: GLF23, Multi-Mode Model, TGLF (not in production version yet)

Transport modules: Gyro-Landau-Fluid (GLF23)



[R E Waltz et al Phys. Fluids B 1992]

- 3D non-linear GLF code (fluctuating n_i , $P_{i/l}$, $P_{i\perp}$, $V_{i/l}$, trapped n_e and P_e , passing n_e , $A_{//}$, etc., quasineutrality condition + 3D ballooning GKS code

- mixing length estimation of transport coefficients:

 $\chi \approx 1.5(\gamma_{net} / k_{xM}^2) \times \gamma_d \gamma / (\gamma^2 + \omega^2)$

- ITG, TEM and ETG turbulence, 10 modes in *k*-spectrum

- dependencies: critical gradient, T_i/T_e , L_{Ti} , L_{Te} , L_n , q

- implementation: GLF23 coefficients computed selfconsistently with Braginskii equations for averaged quantities in time dependent simulations

- computes χ_e , χ_i , D_i , χ_ϕ

 available in NTCC module library: http://w3.pppl.gov/rib/repositories/NTCC/catalog/



Transport modules: Multi-Mode Model



[G Bateman et al Phys. Plasmas 1998]

 $\chi = \chi(centre) + \chi(gradient region) + \chi(edge) + \chi_{neocl.}$

Kinetic Ballooning Model Weiland model Horton ETG model RBM, DRIBM by T. Rafiq (L-mode) NCLASS paleoclassical

- Flexible implementation: various components can be switched on/off
- Detailed output: χi , χe , $\chi \varphi$, D for all models, growth rates for Weiland and DRIBM)

T. Rafiq et al, submitted to PoP



Momentum transport



TRANSP variables and equations relevant to momentum transport: angular velocity:

 $\omega(\sqrt{\Phi}) = V\varphi/R$ (all thermal species are assumed to have the same angular velocity)

angular momentum density (in analysis mode inferred from density, Z_{eff} and toroidal rotation measurements): $m = \langle R^2 \rangle_{\omega} \Sigma_i n_i^* M_i$ (sum over thermal species, electron momentum is ignored)

input torques:

 $S\phi_{total} = S_{nbi,coll}(including S_{nbi,i} \& S_{nbi,e}) + S_{nbi,therm} + S_{JxB}(radial movement of fast ions) + S_{ripple} + S_{anom}(prescribed)$

CX losses:

S_{CX_losses} = S₁(beam CX momentum loss (NUBEAM)) - S₂(partly recaptured by neutral gas model (FRANTIC))



Momentum balance equation:

dm/dt = S ϕ_{total} - S_{CX_losses} - (1/r)(d/dr)($\chi_{\phi}\nabla m$) - <R²>ω $\Sigma_{j}\Gamma_{j}M_{j}$

 $\chi_{\scriptscriptstyle \Phi}$ is the momentum diffusivity calculated in analysis mode

Momentum confinement time :

 $dL/dt = \int (S\phi_{total} - S_{CX_losses} - \langle R^2 \rangle \omega \Sigma_j \Gamma_j M_j) dv - L/\tau_{\phi} \qquad L = \int (\langle R^2 \rangle \omega \Sigma_j n_j^* M_j) dv$

^{*} Description of toroidal momentum transport in TRANSP: R. J. Goldston in "Basic Physical Processes of Toroidal Fusion Plasmas", Varenna, 1985 ftp://ftp.pppl.gov:/pub/dmccune/transp/Goldston_Rotation.pdf

Modelling of trace species

$$\frac{dn}{dt} = \mathsf{S} - \mathsf{div}(\Gamma)$$

Here S includes beam fuelling, gas flow and recycling

$$\Gamma = V * (-Lf *grad(n) + n)$$
 (1)
 or
 $\Gamma = -D *grad(n) + V *n$ (2)

V is the non diffusive radial velocity, Lf is the diffusive flow scale length, D is the specie diffusivity (Ufile input). The following combinations are allowed: V only -- formula (2), assuming D=0.0 D only -- formula (2), assuming V=0.0 D and V -- formula (2) Lf and V -- formula (1)

Evolution of trace tritium density after the short tritium gas puff [Voitsekhovitch et al, PoP 2005]



Summary

- TRANSP has incorporated state-of-the-art source and transport models for use as either an interpretive or predictive analysis tool
 - Driven by community input and needs
- Treatments and solvers are becoming increasingly sophisticated and comprehensive
- Provide links to other computationally intensive codes (macro-, micro-fast ion stability, RF full wave, etc.)
- Extensive use of TRANSP capabilities (presently and planned) by a large number of institutions around the world
- Running TRANSP on FusionGrid allows for timely support and quick debug turnaround by PPPL Computational Plasma Physics Group

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Areas of developments

- Continued PT_SOLVER development
 - Momentum, density predictions, benchmark with TGYRO
- Free boundary equilibrium
 - Shape control algorithms, more flexible current control
- NUBEAM and RF
 - TF ripple, coupling to RF, assess GPU for speedup, upgrade RF modules
- Regression testing for code upgrades
- Parallel architecture
 - NUBEAM/TORIC simultaneous execution, solver speedups
- Output pipeline
 - More flexible and state-of-the-art outputs and graphics
- Update tools for preparing and launching runs, assess UFILE structure
- Additional physics
 - Improved (2D) neutral transport model, physics-based AE-driven fast ion diffusion model, long term (edge-core, stability, turbulence code interfaces)
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